

## “Wet-to-Dry” Conformational Transition of Polymer Layers Grafted to Nanoparticles in Nanocomposite

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Soft polymer properties are improved by inclusion of small hard inorganic particles inside the melt matrix. More recently, special attention has been focused on the reduction in filler size down to the nanometer range to enhance the macroscopic properties of the nanocomposites for the design of new materials for applications in mechanical, optical, fuel cells, or gas barrier engineering. The key point is the compatibility between the nanofiller and the polymer matrix, and one active field is chain grafting synthesis to dress the particle with a polymer corona of the same nature as the polymer matrix. Many recent studies present various refined synthesis processes of controlled polymerization to obtain well-defined grafted nanoparticles.<sup>1</sup> Depending on the synthesis conditions, it is possible in a controlled way to change the grafting density, the length of the grafted chains, or the particle's colloidal stability. The resulting grafted objects can be used as fillers by mixing with free chains of the same polymer to create new hierarchical structures inside the matrix, from well-dispersed state to large anisotropic or fractal aggregates without direct connectivity, or with formation of a continuous network.<sup>2</sup> Theoretical interactions have been described for planar<sup>3,4</sup> or curved surfaces.<sup>5–7</sup> The free energy of swelling can be expressed by two terms: a mixing entropy term between the grafted and the free chains and an elastic term of deformation (stretching or compression) of the grafted chains. The low mass free entities (solvent molecules, monomer, or polymer chains corresponding to a grafted-free chain length ratio superior to one) can penetrate into the grafted corona according to a favorable entropic potential and swell it: the grafted chains are stretched, and the corona is “wet”. Conversely, when grafted and free chain masses are comparable (grafted-free chain length ratio equal or inferior to one), mixing entropy is much lower and swelling elastic energy dominates, which expels free entities from the corona and changes its conformation. As a consequence, the grafted chains are compressed and the corona becomes “dry”. These concepts have been studied experimentally in the case of hydrogels<sup>8,9</sup> or large silica particles<sup>10</sup> through indirect measurements of the interparticles structure factor with X-rays or light scattering.

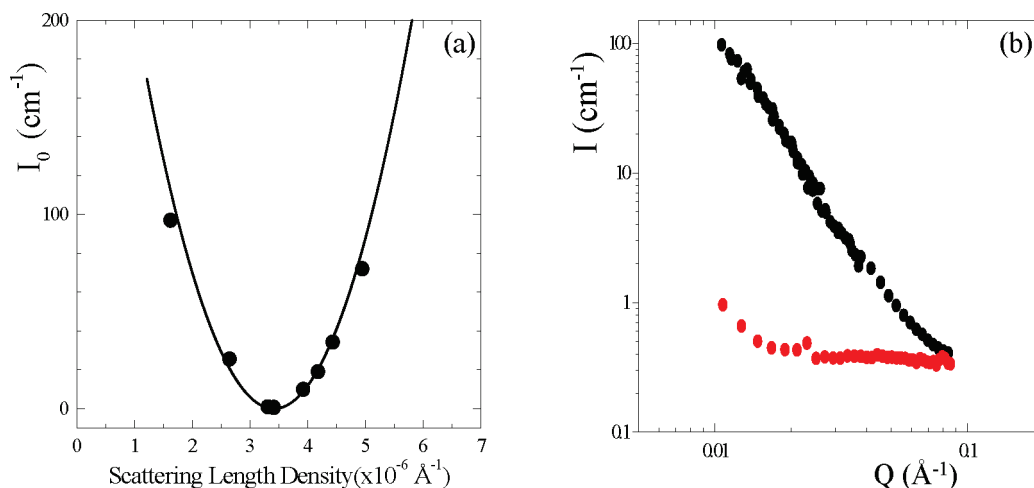
But a direct measurement of the grafted polymer conformation and its change due to specific interactions with free chains is still missing, particularly in a nanocomposite system in which it is technically difficult to distinguish the grafted chains from the free ones.

The present Communication reports the first direct measurement of the conformation of the polymer brush grafted around silica nanoparticles inside a matrix of the same polymer. This success is based on two main contributions. First, the development of refined chemistry to build controlled grafted silica particles with deuterated polystyrene (PS) chains and to synthesize a statistical hydrogenated–deuterated polystyrene matrix which matches the particles scattering in small-angle neutron scattering (SANS) experiments: the silica becomes invisible to neutrons in the matrix, and only the deuterated corona contributes to the signal. Second, a combination of small-angle and ultrasmall-angle X-ray scattering (SAXS/USAXS) to extract the interparticles structure factor, corresponding to the organization of the grafted particles inside the polymer matrix with microscopy (TEM) to check the particle dispersion in real space.

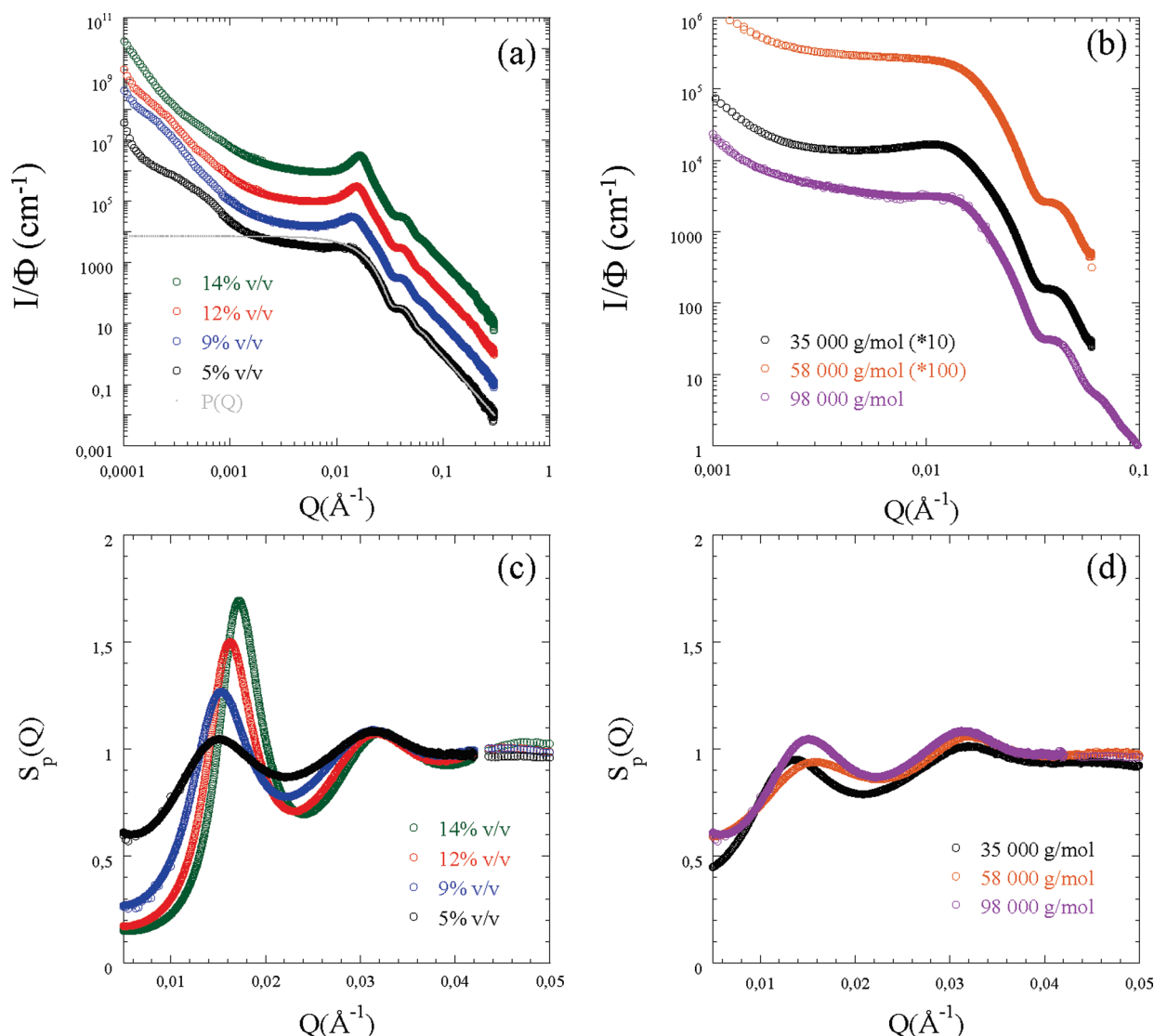
We use a “grafting from” technique associated with a radical-controlled polymerization (nitroxide-mediated polymerization, NMP) to obtain well-defined silica particles (Ludox type) grafted with polystyrene chains. The grafting density is fixed at 0.2 molecules/nm<sup>2</sup>, and the polymerization gives chains of  $M_n = 25\,000$  g/mol with low polydispersity index around 1.2. The complete description of the grafting process can be founded in ref 11. The synthesis process permits to graft either normal or deuterated polymers at the surface of the particles, but only deuterated grafted chains were used in the present paper. The grafted corona can be modeled in solution using SANS contrast variation method with a core–shell model from which we can extract a layer thickness.<sup>11</sup> We have then well-defined grafted nanoparticles, which can be suspended in an organic solvent (dimethylacetamide, DMAc) and mixed with PS chains at 10% v/v dissolved in the same solvent. Final nanocomposites are formed according to a controlled solvent slow evaporation process described previously<sup>12</sup> which permits to obtain homogeneously filled films samples. The polymer chains for the matrix were synthesized by conventional radical polymerization (6 h at 80 °C) by mixing normal and deuterated styrene to form a statistical H/D PS chain at the final state of the polymerization process. The styrene H/D ratio was varied (from 35 to 100% v/v of H-monomer) and the neutron scattering length density (SLD) of the resulting chain measured with SANS (PAXY at LLB) to find the SLD value the closest possible of the one of the silica ( $3.40 \times 10^{-6} \text{ Å}^{-2}$ ) (Figure 1a). To check this, a “blank” filled with bare silica is used: the scattering coming from the silica is reduced by more than 2 orders of magnitude (Figure 1b). Varying the polymerization time allows to synthesize three different masses for the matrix, respectively of  $M_n = 35\,000$ , 58 000, and 98 000 g/mol.

Different nanocomposites were then formed either at fixed grafted/free matrix chain ratio ( $M_{n,\text{grafted}} = 25\,000$  and  $M_{n,\text{matrix}} = 98\,000$  g/mol) with growing quantities of grafted silica particles (from 5 to 20% v/v) or at fixed particles content (5% v/v) by varying the chain length ratio ( $M_{n,\text{grafted}} = 25\,000$  and  $M_{n,\text{matrix}} = 35\,000$ , 58 000, and 98 000 g/mol). The samples were analyzed with a combination of SAXS and USAXS measurements on ID02 at ESRF using the classical pinhole setup and the Bönse-Hart setup. Results are presented in Figure 2a for concentration variation and in Figure 2b for chain length ratio variation. As X-rays do not differentiate between labeled polymers, these spectra reveal the

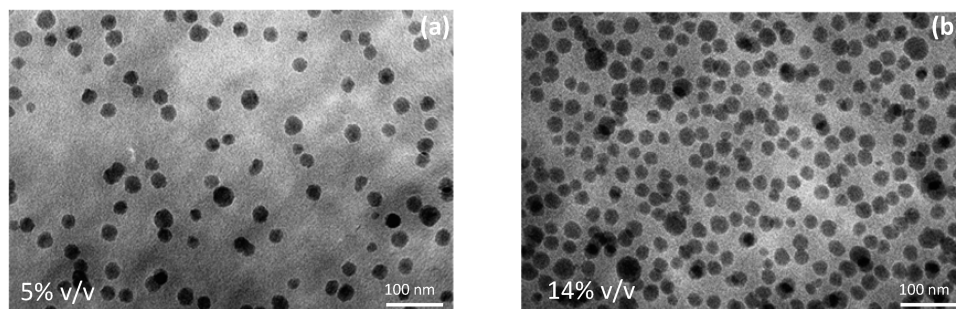
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**Figure 1.** (a) Variation of the H/D styrene ratio for synthesis of a statistical H/D PS matrix ( $M_n = 98\,000$  g/mol) of tunable scattering length densities (SLD): the minimum of the curve correspond to the SLD of the silica. (b) Two nanocomposites filled with nongrafted silica for different SLD matrix values, far from the silica SLD point (black circles) and close to the silica SLD point (red circles).



**Figure 2.** SAXS and USAXS normalized intensities (ID02 at ESRF) on PS nanocomposites filled with deuterated PS grafted silica particles as a function of (a) the silica volume fraction ranging from 5% to 14% v/v at fixed chain length ratio ( $M_{n, \text{grafted}} = 25\,000$  g/mol and  $M_{n, \text{matrix}} = 98\,000$  g/mol) and (b) the chain length ratio at fixed silica volume fraction (5% v/v) for  $M_{n, \text{grafted}} = 25\,000$  g/mol and  $M_{n, \text{matrix}} = 35\,000$ , 58 000, and 98 000 g/mol. The full gray line in (a) represent the polydisperse sphere form factor of the native silica particle  $P_p(Q)$ ; (c) and (d) are the corresponding interparticle structure factor  $S_p(Q)$  inside the polymer matrix obtained by divided the total intensity (a) and (b) by the form factor of a single particle (full gray line).



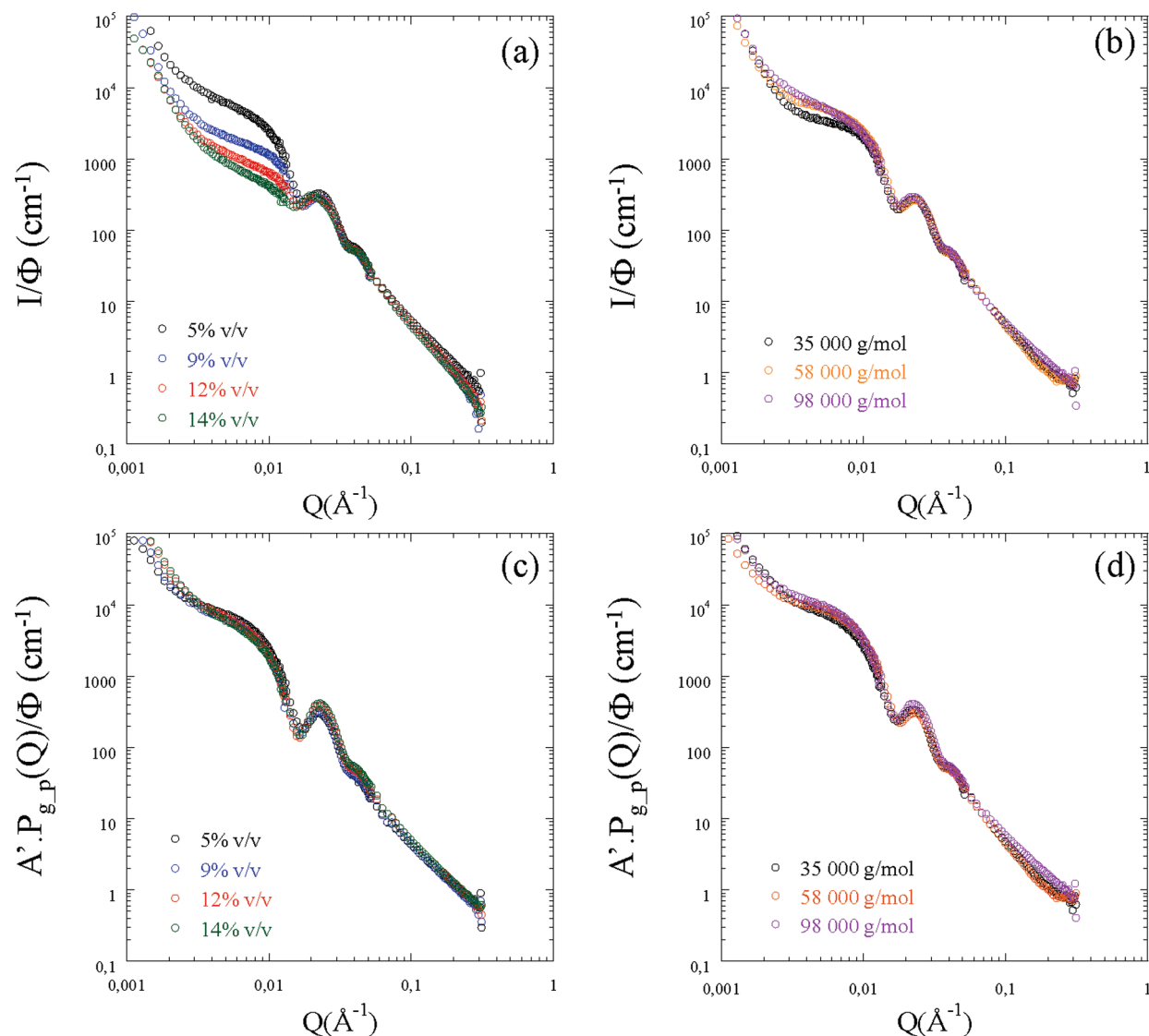
**Figure 3.** Transmission electron microscopy (TEM) on the nanocomposite ( $M_{n,\text{matrix}} = 98\,000$  g/mol) filled with 5% v/v (a) and 14% v/v (b) of grafted silica particles ( $M_{n,\text{grafted}} = 25\,000$  g/mol).

silica particles dispersion inside the polymer matrix related to local organization of the particles between themselves, characterized by the interparticles structure factor  $S(Q)$ . The strong upturn in the very low  $Q$  range comes from the formation of void cavities inside the materials related to the processing formation used for the films.<sup>13</sup> According to the general scattering equation of a particle system in a continuous medium  $I_{\text{SAXS}}(Q) \sim A S_p(Q) P_p(Q)$  ( $A$  is proportional to the product of the contrast term  $\Delta\rho^2$ , the volume of the scattered particle and the number of scattered objects; the index  $p$  represent the silica particles), the interparticles structure factor can be calculated by dividing the total scattered intensity,  $I(Q)$ , by the single form factor of the silica particle  $P(Q)$ . The form factor of the silica particles is well-known in our case as it can be determined easily through the analysis of the scattering of a diluted solution of particles (see the full gray line in Figure 2a). The silica particles form factor follows the classical behavior of polydisperse (log-normal) spheres of a mean radius of  $\langle R \rangle = 13$  nm with a dispersion of 0.14. The interparticles structure factors deduced from the division of the intensity by the silica form factor are presented in Figure 2c for the evolution as a function of the particle concentration and in Figure 2d for the evolution as a function of the chain length ratio. The appearance of a strong correlation peak indicates a well-dispersed repulsive organization of the particle inside the polymer matrix with a first neighboring typical distance equal to  $2\pi/Q^*$  ( $Q^*$  is the position of the first maximum peak). The evolution of the peak position  $Q^*$  with the silica concentration (Figure 2c) follows the classical behavior of a 3D cubic volume system while the peak position does not move as a function of the chain length ratio (Figure 2d), indicating that the dispersion of the particle does not change in the probed chain length ratio range. The good dispersion of the particles is confirmed in real space by the microscopy pictures presented in Figure 3 for two particle concentrations (respectively 5% v/v Figure 3a and 14% v/v Figure 3b).

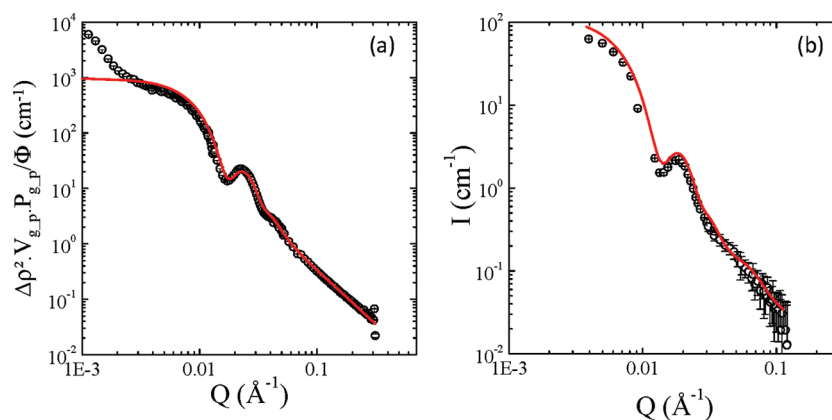
The same samples were then measured by SANS (D11 at ILL), and the scattering curves are presented in Figure 4a as a function of the grafted silica concentration and (b) as a function of the chain length ratio. The scattered intensity is thus only due to the grafted deuterated chains around the silica particles as the silica signal contribution is suppressed by the matrix one and can be expressed in this case as  $I_{\text{SANS}}(Q) \sim A' S_{\text{gp}}(Q) P_{\text{gp}}(Q)$ , where the index "gp" represents the grafted particles. For a correct modeling of the grafted brush ( $P_{\text{gp}}(Q)$ ), it is needed to suppress the contribution of the interparticles structure factor ( $S_{\text{gp}}(Q)$ ) related to the arrangement of the grafted coronas between them inside the matrix. As this organization is built across the network center of masses of the scattered grafted objects, the interparticles structure factor is independent of the nature of the probe, i.e., is the same for neutron and for X-rays,  $S_p(Q) = S_{\text{gp}}(Q)$ . Thus, the grafted corona form factor can be extracted by dividing the SANS signal (Figure 4a,b) by the particles structure factor deduced from SAXS (Figure 2c,d). The results are presented in

Figure 4c for the evolution as a function of the particles content and Figure 4d for the evolution as a function of the chain length ratio. Figure 4c shows a nice superimposition of the different curves, illustrating that the form factor of the grafted corona does not depend on the silica concentration. In other words, increasing the grafted interparticles interactions by decreasing the mean distance between themselves does not change the conformation of the grafted brushes. In addition, Figure 4d shows also a nice superimposition of the different curves. Changing the interaction between the grafted and the free matrix chains is expected to have an influence on the brush conformation: reducing the free chain masses from 98 000 to 35 000 g/mol was supposed to reduce the mixing entropy term and induce an increase of the stretching rate of the grafted brushes. The grafted corona form factor can be perfectly modeled (Figure 5a) with a Gaussian chain model<sup>14</sup> from which we can extract the physical parameters describing the system: the number of grafted silica particles (5% v/v), the number of polymer chains per particle ( $\sim 400$ ), the mass ( $M_n = 25\,000$  g/mol), and the radius of gyration of grafted chains equal to 3.1 nm, corresponding to a corona thickness of 6.2 nm. The conformation of the grafted corona inside the composite is compared with the conformation of the same grafted corona in solution, before the mixing with the matrix chains. For such a solution, matching the silica signal is easy using the appropriate fractions of deuterated and hydrogenated solvent. The result, presented in Figure 5b already published in ref 11, is perfectly fitted by a core-shell model with a homogeneous corona of thickness of 12 nm.

We thus observe a clear-cut change of the grafted corona conformation from the solution to the nanocomposites: the grafted chains are stretched in solution while in the nanocomposites they return to a less expanded Gaussian coil conformation. The observation is in very good agreement with the theoretically expected wet-to-dry transition of the grafted corona conformation induced by the mixing entropy potential between the free and the grafted chains. In addition, we demonstrate that interactions with longer free chains induce a compression of the grafted chains of a factor 2 in thickness to minimize the corresponding free energy. No effect of the interparticles distance or of the variation of the grafted-free chain length ratio in the composite has been observed on the amplitude of the elastic compression of the corona. But such effect is expected to be observed for larger parameters: for higher silica volume fraction or large differences in chain length ratios. This measurement and modeling of the polymer corona grafted on particles inside a polymer matrix is thus an experimental breakthrough thanks to a combination of pertinent chemical synthesis and of complementary scattering techniques SANS and SAXS and is important on two levels: it validates the theoretical predictions, and it enables better understanding of the dispersion mechanisms of nanoparticles in nanocomposites and specifically the contribution of the brush conformation modifications in these mechanisms.



**Figure 4.** (a) SANS curves of the deuterated grafted corona ( $M_{n, \text{grafted}} = 25\,000$  g/mol) inside the silica matched matrix ( $M_{n, \text{matrix}} = 98\,000$  g/mol) at four particle concentrations (5, 9, 12, and 14% v/v). (b) SANS curves of the deuterated grafted corona ( $M_{n, \text{grafted}} = 25\,000$  g/mol) at silica volume fraction of 5% v/v inside the three silica matched matrix ( $M_{n, \text{matrix}} = 35\,000$ , 58 000, and 98 000 g/mol). (c) Evolution of the form factor of the grafted polymer corona in the nanocomposite  $P_{g,p}(Q)$  as a function of the silica content obtained after division of the neutron scattering intensity (a) by the interparticles structure factor  $S_p(Q)$  (Figure 2c). (d) Evolution of the form factor of the grafted polymer corona in the nanocomposite  $P_{g,p}(Q)$  as a function of the chain length ratio obtained after division of the neutron scattering intensity (b) by the interparticles structure factor  $S_p(Q)$  (Figure 2d).



**Figure 5.** (a) Modelization of the form factor of the deuterated PS grafted silica nanoparticles ( $M_{n, \text{grafted}} = 25\,000$  g/mol) dispersed in a silica matched matrix ( $M_{n, \text{matrix}} = 98\,000$  g/mol) with a Gaussian chain model for comparison with the modelization of the same grafted particles dispersed in a solvent (b) and modeled with a core-shell model (from ref 11).

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